PHYSICOCHEMICAL PROPERTIES OF Sn(S_{1-x}Te_x) SOLID SOLUTIONS OF BOTH MASSIVE MATERIALS AND THIN FILMS

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The $Sn(S_{1-x}Te_x)$ solid solutions have been prepared on both massive materials and thin films with variation of x value of 0, 0.2, 0.4, 0.6, 0.8, and 1.0. The massive materials were prepared by Bridgman technique, while the thin films were obtained with vacuum evaporation technique. Phase and crystal structure of both massive materials and thin films were investigated by X-ray diffraction (XRD). The chemical compositions were determined by energy dispersive X-ray spectroscopy (EDAX). Especially for thin films, the surface morphology was studied by scanning electron microscopy (SEM), while the optical properties were obtained by UV-Vis-NIR spectrophotometers. The results show that the crystalline structures of all materials have an orthorhombic structure when dominated by sulfur atom, while having a cubic structure for tellurium atom dominance. All samples are p-type semiconductors and the optical properties show a reduction of the band gaps when the tellurium atomic composition increases.

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1. Introduction

Tin chalcogenide nanoparticles, SnS and SnTe have received considerable attention because of their interesting optoelectronic properties. Tin monosulfide (SnS) is p-type semiconductors that attract the attention of the researchers in the context of photovoltaic conversion. Bulk SnS has both a direct band gap ($E_{g,dir} = 1.32 \text{ eV}$) and indirect band gap ($E_{g,ind} =$ 1.08 eV) [1, 2], which is very close to that of silicon and is especially attractive due to its high conversion efficiency in photovoltaic devices $(\sim 25\%)$ [3]. The efforts to produce new photovoltaic materials with high efficiency are needed because each country will experience a conventional energy crisis as well as the pollution issue it generates. Crystal of SnS crystallizes in the orthorhombic structure of space group *Pbnm* [4, 5]. The SnS is of crystallographic interest because it is intermediate between three-dimensional crystalline networks and two-dimensional layer type compounds [6]. Tin telluride (SnTe) is an important narrow band gap semiconductor material. Bulk SnTe is a semiconductor with a direct band gap of 0.35 eV [7-9], which is narrower than that of the bulk SnS. Kovalenko et al. (2007) report a solution-phase synthesis of high-quality colloidal SnTe nanocrystals with mean diameters tunable in the range of ca. 4.5-15 nm and corresponding band gaps of 0.80- 0.38 eV [10]. At room temperature, crystal of SnTe crystallizes in the cubic structure of space group *Fm-3m* [11].

The doping tellurium (Te) on SnS is expected to reduce band gap of SnS so that obtained material with band gap in accordance with maximum theoretical conversion efficiency in solar cell [12]. In this paper, we report the preparation of $Sn(S_{1-x}Te_x)$ bulks by Bridgman technique and $Sn(S_{1-x}Te_x)$ thin films by vacuum evaporation techniques. Therefore, this study specifically aims to determine the effect of atom fraction of Te on SnS. The effect is estimated on two important points. First, it relates in the crystal structure (lattice parameter). This point corresponds to the

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dependence of the $Sn(S_{1-x}Te_x)$ lattice parameter to the tellurium atom x fraction. Secondly, it relates to the band gap width dependence as a function of the fraction x of the tellurium atom. These two physical quantities are very important in relation to the p-n junction of solar cells, especially in the ability of material absorption to solar energy. The optical and structural properties of the $Sn(S_{1-x}Te_x)$ solid solutions have been investigated on both massive materials and thin films. Its were studied using UV-Vis-NIR spectrophotometer, X-ray diffraction (XRD), scanning electron microscopy (SEM), and energy dispersive of X-ray spectroscopy (EDAX).

2. Experimental

2.1. Sample Preparation

The materials needed for the preparation of massive materials of solid solution are tin (Sn), sulfur (S) and tellurium (Te). Each material has a degree of purity 99.99%. The study was distinguished in three steps. First step, preparation of bulks of $Sn(S_{1-x}Te_x)$ solid solutions using Bridgman technique. For the preparation of $Sn(S_{1-x}Te_x)$, first weighed tin (Sn) eg. p gram. Further calculated sulfur mass S using the following formula: $\left(\frac{p}{(BA)_{Sn}}, (1-x), (BA)_S\right)$ gram, while the mass of

tellurium (Te) can be calculated by the formula: $\left(\frac{p}{(BA)_{Sn}}(x)(BA)_{Te}\right)$ gram, where, BA denotes the atomic

weight. Furthermore, the three metal materials: Sn, S and Te are inserted in pyrex tubes having an inner and outer diameter of 12 and 16 mm, respectively. The tubes were previously washed with a mixture of HF, HNO_3 and H_2O solutions with a ratio of 2: 3: 5 and dried in a heating chamber at 80°C for 8 hours. The tube containing the above materials was placed on a 10⁻⁵ Torr vacuum and welded at one end, resulting in a capsule containing: Sn, S and Te. The capsule was then placed in a furnace whose temperature and duration of maintaining the temperature can be adjusted as required in the form of a heating flow. The second step, thin films were reproduced using vacuum evaporation techniques. The evaporator reactor system was made of 30 cm diameter quartz tube placed over a vacuum system comprising a primary pump and a secondary pump. As the source or target was a solid solution of $Sn(S_{1,x}Te_x)$ prepared by Bridgman technique which was then crushed so as to form powder. The powder was placed on a boat-based heating source that can be heated to a temperature above 1500°C made of molybdenum (Mo) material. The temperature of the cup was measured using a thermocouple. Between source and substrate separated by a separator whose distance can be varied at 10, 15, and 25 cm. Meanwhile around the substrate was installed substrate heating due to crystallization of the thin layer will occur with substrate temperature above 300°C [13]. The advantage of this technique was that the thin film yields the same composition and structure as the source of the massive solid solution target and the working pressure of 10⁻⁵ Torr, thus using primary and secondary pumps or diffusion pumps.

2.2. Sample Characterization

The chemical compositions of massive materials and thin films of $Sn(S_{1-x}Te_x)$ were determined by energy dispersive X-ray spectroscopy (EDAX) from EDAX Inc., while the surface morphology of thin films were investigated by scanning electron microscopy (SEM) (Jeol JSM-6360). The machine was operated at an acceleration voltage of 20 keV at a working distance of 10 mm to identify the morphological properties of samples.

The prepared samples were examined for the investigation of phase and structure using a Rigaku Miniflex 600-benchtop X-ray Diffraction (XRD), operated in the Bragg configuration using Cu K α radiation ($\lambda = 1.5406$ Å). The intensities were determined in the 2 θ collection range of 20° to 90°. The XRD instrument was run at 40 kV and 15 mA. The lattice parameters and Bravais lattice from the powder diffraction data were carried out using least squares refinement method by the U-fit program [14].

Ultra violet-visible (UV-Vis) absorption measurements of thin film samples were performed by UV-Vis spectrophotometers (Shimadzu, UV-2600) in the wavelength range of 200-900 nm. The absorption coefficient α on UV-VIS spectroscopy satisfies the following equation [15]:

$$\alpha$$
. $hv = B (hv - E_g)^n$

where h is Planck's constant $(6.626 \times 10^{-34} Js)$, v is the wave frequency (Hz), B is a constant of 10^5 to 10^6 , while n is $\frac{1}{2}$ for direct band gap, where n is 1 for indirect transitions, hv is the energy of photons in electron volt (eV), E_g is the band gap or energy gap (eV). The band gap can be determined by performing absorbance data analysis that is the relationship with absorbance coefficient A with the absorption coefficient α . The relationship of both is linear, whereas in semiconductor materials apply the equation as mentioned above with $n = \frac{1}{2}$. The band gap value (E_g) can be determined which is the value when the left side is equal to zero:

$$(\alpha. h\nu)^2 = B^2 (h\nu - E_g)$$

3. Results and discussion

3.1. SEM and EDAX

EDAX characterization is performed on both a massive sample and a thin film. Table 1 shows the chemical composition of massive materials of $Sn(S_{1-x}Te_x)$ solid solutions and the thin film as shown in the Table 2. The EDAX results show the compatibility of chemical composition between massive material and thin film. This shows the superiority of vacuum evaporation techniques.

Solid Solutions	Chemi	Solid Solutions		
Expected	Sn (%)	S (%)	Te (%)	Obtained
SnS	52.79	47.15	-	SnS _{0.89}
$Sn(S_{0.80}Te_{0.20})$	53.27	37.67	9.06	$Sn(S_{0.71}Te_{0.17})$
$Sn(S_{0.60}Te_{0.40})$	50.26	28.26	20.88	$Sn(S_{0.56}Te_{0.44})$
$Sn(S_{0.40}Te_{0.60})$	50.86	11.94	37.19	$Sn(S_{0.20}Te_{0.70})$
$Sn(S_{0.20}Te_{0.80})$	46.86	9.08	44.06	$Sn(S_{0.19}Te_{0.94})$
SnTe	53.6	-	46.95	SnTe _{0.88}

Table 1. Chemical composition of of $Sn(S_{1-x}Te_x)$ massive materials

Solid Solutions	Chemi	Solid Solutions		
Expected	Sn (%)	S (%)	Te (%)	Obtained
SnS	55.00	45.00	-	SnS _{0.80}
$Sn(S_{0.80}Te_{0.20})$	51.20	40.41	8.39	$Sn(S_{0.79}Te_{0.16})$
$Sn(S_{0.60}Te_{0.40})$	51.87	28.41	19.72	$Sn(S_{0.55}Te_{0.38})$
$Sn(S_{0.40}Te_{0.60})$	54.54	10.85	34.61	$Sn(S_{0.18}Te_{0.64})$
$Sn(S_{0.20}Te_{0.80})$	52.31	0.88	46.81	$Sn(S_{0.02}Te_{0.89})$
SnTe	51.92	-	49.08	SnTe _{0.94}

Table 2. Chemical compositions of $Sn(S_{1-x}Te_x)$ thin films

The surface morphology of thin films was investigated by SEM. Figure 1 show SEM images of $Sn(S_{1-x}Te_x)$ thin films. SEM images show that the surface structures of thin films with sulfur and tellurium dominant are significantly different. The first surface of Fig. 1(a), 1(b), and 1(c) related to the orthorhombic structures. The second surface of Fig. 1(d), 1(e) and 1(f) looks related to the cubic structures.



e) $Sn(S_{0.02}Te_{0.89})$



Fig. 1. SEM images of $Sn(S_{1-x}Te_x)$ thin films

3.2. Structural Characterization

Fig. 2 and 3 depict the XRD patterns of $Sn(S_{1-x}Te_x)$ solid solutions of massive materials and thin films, respectively. The XRD patterns were recorded in 20 range of 15-90°. These diffraction peaks are consistent with the reflections of SnS phase and SnTe. The XRD results of tin monosulfide (SnS) and tin monoteluride (SnTe) are completely different characteristics.



Fig. 2. XRD patterns of $Sn(S_{1-x}Te_x)$ massive materials



Fig. 3. XRD patterns of $Sn(S_{1-x}Te_x)$ thin films

X-ray diffraction pattern shows the peaks that appear in the sample. XRD results data obtained from the research sample. The crystal lattice parameters (a, b and c) formed can be calculated by least squares refinement method [14]. The calculation results show that the crystal structure of SnS is orthorhombic, while the SnTe is cubic with lattice parameters as in Table 3. The massive materials and thin films of $Sn(S_{1-x}Te_x)$ solid solutions have the same structure. The crystal structure in the thin films is orthorhombic when the dominance of sulfur atoms, and will becomes cubic structure when the dominance of tellurium atoms. X-ray diffraction peaks of SnS represent orthorhombic structures with space group *P*bnm, whereas SnTe is related to cubic structures with space group *F*m-3m. The peaks of the X-ray diffractogram emerging from the SnS and SnTe crystals relate to crystalline planes (hkl) or Miller indices. For SnS crystals, the emerging peaks are as diverse as peaks (111) corresponding to 2θ at 31.8° , while for SnTe crystals, the highest peak at an angle of 29.71° indicates the crystalline planes (200). Overall, the peaks appearing on the SnTe crystals represent the cubic structure with the bravais lattice of the center cubic. The result of calculation of lattice parameters of each Sn(S_{1-x}Te_x) solid solutions of both massive materials and thin films can be given in Table 3.

$Sn(S_{1-x}Te_x)$ Solid	Massive Materials			Crevetel System	Figure of Merite			
Solutions	a (Å)	b (Å)	c (Å)	Crystar System	D	R		
SnS	4.317	11.447	3.895	Orthorombic, Pbnm	0.022	0.025		
$Sn(S_{0.71}Te_{0.17})$	4.329	11.196	3.984	Orthorombic, Pbnm (major)	0.032	0.051		
$Sn(S_{0.56}Te_{0.44})$	4.337	11.205	3.991	Orthorombic, Pbnm (major)	0.027	0.039		
$Sn(S_{0.20}Te_{0.70})$	6.319	6.319	6.319	Cubic, Fm-3m (major)	0.009	0.030		
$Sn(S_{0.19}Te_{0.94})$	6.320	6.320	6.320	Cubic, Fm-3m (major)	0.013	0.027		
SnTe	6.322	6.322	6.322	Cubic, Fm-3m	0.042	0.056		
	Thin Films							
SnS _{0.80}	4.293	11.217	3.910	Orthorombic, <i>P</i> bnm	0.028	0.031		
$Sn(S_{0.79}Te_{0.16})$	4.286	11.118	3.946	Orthorombic, Pbnm (major)	0.021	0.031		
$Sn(S_{0.55}Te_{0.38})$	4.230	11.190	4.030	Orthorombic, Pbnm (major)	0.042	0.053		
$Sn(S_{0.18}Te_{0.64})$	6.303	6.303	6.303	Cubic, Fm-3m (major)	0.031	0.068		
$Sn(S_{0.02}Te_{0.89})$	6.304	6.304	6.304	Cubic, Fm-3m (major)	0.044	0.062		
SnTe _{0.94}	6.309	6.309	6.309	Cubic, Fm-3m	0.063	0.092		
The characteristic factors of the refinement (D) are the mean deviation between $2\theta_{obs}$ and $2\theta_{calc}$, that is								
$D = \frac{1}{n_{hkl}} \sum d(2\theta) $ and the confidence factor (R) given by the following relation: $R = \frac{1}{n_{hkl} - n_{var}} \sum (2\theta_{obs} - 2\theta_{calc})^2$,								
where n_{hkl} is the number of reflections taken into account and n_{var} is the number of refined variables								

Table 3. Lattice parameters of both massive materials and thin films of $Sn(S_{1-x}Te_x)$ solid solutions

3.3. Optical Properties

The result of UV-VIS in absorbance form as a photon wave function can be shown in Figure 4. Figure 4 shows that the absorbance is highly dependent on the material band gap. When the photon comes with the photon energy smaller than the band gap material then there is almost no photon uptake by the material. Furthermore, when the photon energy is the same as the band gap then there is a very drastic change that is the appearance of photon uptake by the material. A special case in this sample is that the photon energy range between the wavelengths of 200 nm to 800 nm or the energy range from 1.5 eV to 6.2 eV.



Fig. 4. Absorbance on a thin layer of (a). SnS, (b). $Sn(S_{0.8}Te_{0.2})$ and (c). $Sn(S_{0.6}Te_{0.4})$



Fig. 5. Graph $(\alpha hv)^2$ as a function of photon energy hv of $Sn(S_{1-x}Te_x)$

The first SnS has a 1.35 eV band gap, while SnTe has a 0.4 eV band gap, so it is believed that Sn(S,Te) has a band gap between 1.35 eV and 0.3 eV. The determination can be shown in the Figure 5. According to Figure 5, the energy gap SnS is obtained at 1.35 eV and decreases to 1.25 eV when there is a tellurium fraction of 20%. The band gap decreases again when the atomic fraction of Te gets larger by 1.09 eV for Te fraction to 40%.

4. Conclusions

The results of the research show the crystal structure and lattice parameters (a, b and c) of massive materials and thin films of the $Sn(S_{1-x}Te_x)$ depends on the dominance of the constituent atoms both sulfur and tellurium atoms. The situation is generally applicable to both massive materials and thin films. If the majority of the constituents are Te atoms then the solid solutions and thin films have a cubic structure whereas if dominance by sulfur atoms then it has an orthorhombic crystal structure. The measurement of the chemical composition of the entire sample preparation is non-stoichiometry which presents a deviation from the hope atomic composition, but in reality the third phase of the atom is already formed in both a massive solid solution and a thin film. The surface morphology shows the grain size indicating the formation of crystals in each of the preparation results on massive material and thin film according to the reality of XRD results. Measurement of all electric properties of $Sn(S_{1-x}Te_x)$ has a p-type conductivity with

resistance in accordance with semiconductor material. The band gap decreases as Te atomic fraction increases.

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References

- [1] W. Albers, C, Haas, H.J. Vink, H. J. D. Wasscher, J. Appl. Phys. 32, 2220 (1961).
- [2] O. E. Ogah, G. Zoppi, I. Forbes, R. W. Miles, 23rd European Photovoltaics Solar Energy Conference, p.2580 (Sept. 2008).
- [3] J. P. Singh, R. K. Bedi, Thin Solid Films 199, 9 (1991).
- [4] B. B. Nariya, A. K. Dasadia, M. K. Bhayani, A. J. Patel, A. R. Jani, Chalcogenide Letters 6(10), 549 (2009).
- [5] P. S. del Bucchia, J. C. Jumas, M. Maurin, Acta Crystallogr. B37, 1903 (1981).
- [6] Logothetidid, S.; Polatoglou, H. M. Phys. Rev. B 36, 7491 (1987).
- [7] A. Ishida, T. Tsuchiya, Takaoka, AIP Conf. Proc. 1399, 127 (2011).
- [8] A. Aggarwal, P. D. Patel, D. Lakshminarayana, J. Crys. Growth 142, 344 (1994).
- [9] R Saini, Pallavi, M. Singh, R. Kumar, G. Jain, Chalcogenide Letters 7(3), 197 (2010).
- [10] M. V. Kovalenko, W. Heiss, E. V. Shevchenko, J. S. Lee, H. Schwinghammer, A. P. Alivisatos, D. V. Talapin, J. Am. Chem. Soc. **129**, 11354 (2007).
- [11] Landolt-Bornstein, Springer-Verlag: Berlin, Vol. 17 (1983).
- [12] A. Goetzberger, C. Hebling, Solar Energy Materials and Solar Cells 62, 1 (2000).
- [13] A. Zouaoui, M. Lachab, M. L. Hidalgo, A. Chaffa, C. Llinares, N. Kesri, Thin Solid Films 339, 10 (1999).
- [14] M. Evain, U-FIT: A cell parameter refinement program, IMN Nantes, France, (1992).
- [15] H. Sakata, H. Ogawa, Solar Energy Materials and Solar Cells 63, 259 (2000).